

On Aerosol and Molecular Extinction of Solar Radiation in Transparency “Window” $\nu = 750 - 1250 \text{ cm}^{-1}$ from Atmospheric Measurements

K. A. Shukurov, A. Kh. Shukurov, and G. S. Golitsyn
 Oboukhov Institute of Atmospheric Physics
 Russian Academy of Sciences
 Moscow, Russia

Introduction

In investigations of variations of optical depth of the vertical column of atmosphere $\tau = -\ln P$, where P = transmittance of the atmosphere, in intervals $\Delta\nu_i$ of transparency “microwindows” i in range of wavenumbers $\nu = 750 - 1250 \text{ cm}^{-1}$, one of the most difficult problems is to separate the contributions of an aerosol and water vapor in τ (spectral position of “microwindows” i (see Figure 1 under numbers $i = 1 - 22$). Optical depth $\tau(m, w)$, measured at air mass m and amount of the precipitable water vapor H_2O in the vertical column of atmosphere w (cm), is possible to express by the sum (Bignell, Saiedy, and Shepard 1963):

$$\tau(m) = \tau_a(m) + \tau_{mol}(m) + \tau_{H_2O}^k(m) + \tau_{H_2O}^c(m) \quad (1)$$

where:

$\tau_A(m)$ = aerosol optical depth

$P_A(m) = \exp[-\tau_A(m)]$ - transmittance of an aerosol

τ_{mol} = optical depth of gases N_2O , CO_2 , CH_4 , O_3 (the contribution in $\Delta\nu_i$ of other gases can be neglected)

$P_{mol}(m) = \exp[-\tau_{mol}(m)]$

$\tau_{H_2O}^k(m) = k_{H_2O}^k \times w \times m$

$\tau_{H_2O}^c(m) = k_{H_2O}^c \times w \times m$ - continuum and selective components of a water vapor optical depth

$k_{H_2O}^k$ and $k_{H_2O}^c$ = corresponding coefficients of the absorption H_2O (cm^{-1})

$P_{H_2O}^k(m) = \exp[-\tau_{H_2O}^k(m)]$

$P_{H_2O}^c(m) = \exp[-\tau_{H_2O}^c(m)]$.

At present, the laboratory investigations have not revealed the nature of variations of dependence $\tau_{H_2O}^k$ upon the water vapor partial pressure e (mb) and air temperature t ($^{\circ}\text{C}$)

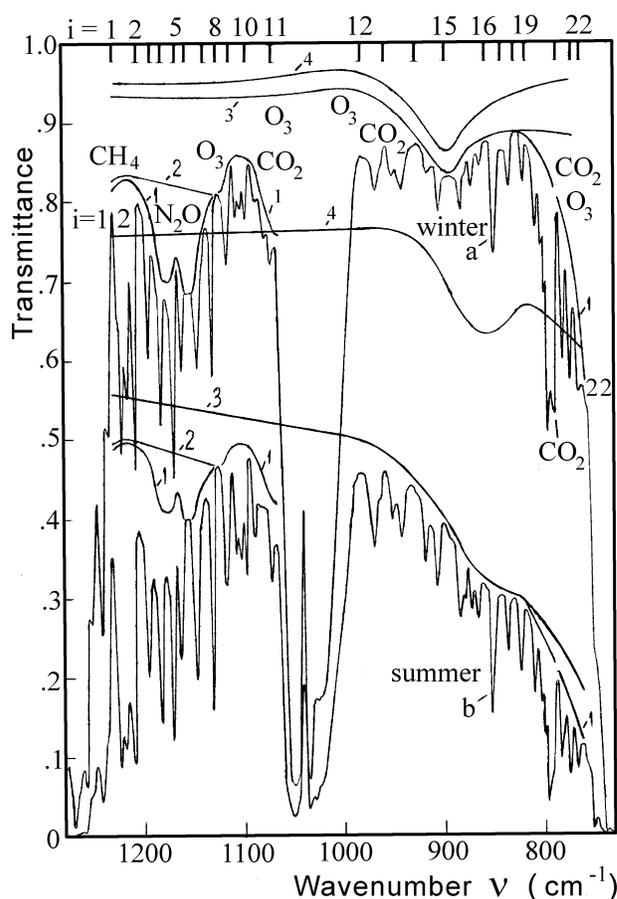


Figure 1. Examples of transmittance P_ν at $m \approx 3$: a - winter, b - summer.

(Kulp and Shinn 1994). This is the main difficulty of an estimation of τ_A , which is necessary, particularly, for calculations of the climatic effects of aerosol and clouds. Another difficulty is connected with the application of the Bouguer's law for determination $\tau_{H_2O}^k$, $k_{H_2O}^k$, and their spectral dependencies (Paramonova, Brounstein, and Frolov

1978). The values τ_i in “microwindows” i are estimated, usually, by Bouguer’s law:

$$\ln I_i = \ln I_i^0 - \tau_i \times m_i \quad (2)$$

where I_i and I_i^0 = measured and top-of-the-atmosphere values of relative intensity I_v of solar radiation in Δv_i .

Values I were measured for estimations of $\tau_{i,A}$ at optimum spectral resolution $\delta v = 1 - 5 \text{ cm}^{-1} \approx \Delta v_i$. Using (2) in Δv_i , where $\tau_{\text{mol}} \neq 0$ and $\tau_{\text{H}_2\text{O}}^c \neq 0$, by calculations of Paramonova et al. (1978) leads to underestimation of I_i^0 and τ_i , as a result of $\tau_i(m) \neq \tau_i \times m$. It was not experimentally checked, since usually I_v were measured at $w \geq 0.5 \text{ cm}$, which didn’t permit the separation of molecular bands, while at small values of w ($\leq 0.3 \text{ cm}$) they are discovered quite clearly (compare spectra P_v on Figure 1).

The goal of the present study is to obtain distinctions $P_{v,A}$ in winter and summer conditions, highly different by w .

Data and Method

The spectra of $P_v = I_v/I_v^0$, obtained in different seasons of 1971 - 1973 at Zvenigorod Scientific Station of IAPh RAS, were used for the estimations of τ and its components. I_v^0 was calculated by interpolation of values I_i^0 obtained from (2) with subsequent adjustment by the recommendations of Paramonova et al. (1978). Spectra of I_v were registered with a two-beam prism spectrophotometer UR - 20 with $\delta v = 4.5 \text{ cm}^{-1}$ and a scanning rate of $40 \text{ cm}^{-1} / \text{min}$. Over 250 realizations I_v were recorded at variations w within the range from $\approx 0.2 \text{ cm}$ to $\approx 2.6 \text{ cm}$ and at near-surface temperature t^0 from $- 25 \text{ }^\circ\text{C}$ up to $+ 30 \text{ }^\circ\text{C}$. Over 100 of them corresponded to $w \leq 0.3 \text{ cm}$ and $t^0 < 0$. Calibration dependencies of P_{min} ($m \times w$)—transmittance in maximum of 2724 cm^{-1} HDO absorption band related to transparency “microwindows” in the region of its wings—were used for estimation of w , which reduced the error of the calculation of w up to $\pm 5\%$. The method of the estimation of components of τ in typical winter and summer conditions is shown in Figure 1 for

$$P_v = P_{v,\text{H}_2\text{O}}^c \times P_{v,\text{mol}} \times P_{v,\text{H}_2\text{O}}^k \times P_{v,A}$$

As an example: a - winter, $w = 0.23 \text{ cm}$, $t^0 \approx - 10 \text{ }^\circ\text{C}$, $m = 3.27$ ($i = 1$) - 3.46 ($i = 22$); b - summer, $w = 2.1 \text{ cm}$, $t^0 \approx + 25 \text{ }^\circ\text{C}$, $m = 3.11 - 3.36$.

The lines 1 - 4 above P_v - (a) were plotted according to estimates of the contribution of gas components (above P_v - (b) they are plotted similarly). The determination of factors P_v were obtained by the following way. The line 1 is $P / P_{\text{H}_2\text{O}}^c$. Values $k_{i,\text{H}_2\text{O}}^c$ in $i = 1 - 5, 7, 9 - 11$ were estimated by dependence of $c_{i,6} = P_i / P_6$ upon m in P_v - (a) and P_v - (b), $k_{i,\text{H}_2\text{O}}^c \approx 0$ was assumed for other i . Since $k_{6,\text{H}_2\text{O}}^c \approx 0$,

$$k_{i,\text{H}_2\text{O}}^c = \Delta k_{i,\text{H}_2\text{O}}^k + [1/(\Delta w \times m)] \times \ln [c_{i,6}(w_a) / c_{i,6}(w_b)] \quad (3)$$

where $\Delta k_{i,\text{H}_2\text{O}}^k = k_{i,\text{H}_2\text{O}}^k - k_{6,\text{H}_2\text{O}}^k$ was calculated from spectral dependence of $k_{i,\text{H}_2\text{O}}^k$ (see below), $\Delta w = 2.1 \text{ cm} - 0.23 \text{ cm} \approx 1.9 \text{ cm}$. A line 2 - $P_{\text{N}_2\text{O}}^0$ (overline of N_2O bands) was plotted by P in $i = 8$ and by values $P / (P_{\text{H}_2\text{O}}^c \times P_{\text{N}_2\text{O}}^0)$ in $i = 1, 2$, where $\tau_{\text{N}_2\text{O}} \approx 0.005 \text{ atm}^{-1}$ by published data. A line 3 - $P^k = P_A \times P_{\text{H}_2\text{O}}^k = P / (P_{\text{H}_2\text{O}}^c \times P_{\text{mol}})$ was drawn by interpolation of values P^k in $i = 2, 8, 9, 12, 15, 16, 18$ and 22 (in $i = 15, 16, 18$ $P^k \approx P$). $P_8^k = P_8 / P_{8,\text{O}_3}$, $P_9^k = P_9 / (P_{9,\text{O}_3} \times P_{9,\text{H}_2\text{O}}^c)$. Values of P_{O_3} for $i = 8$ have been estimated at O_3 content of 0.3 atm.cm according to Bartman et al.(1975). In $i = 2$ value $P_{\text{H}_2\text{O}}^c \times P_{\text{mol}}$ was determined by dependence of P_2 / P_8 upon m . In $i = 12$ $P_{\text{H}_2\text{O}}^c / P_{\text{mol}} = P_{9,\text{H}_2\text{O}}^c \times P_{9,\text{mol}}$. In $i = 22$ $P_{22,\text{mol}} \times P_{22}^c$ was calculated by dependence of P_{22} / P_{18} upon m with correction for Δk^k and with assumption $P_{18,\text{H}_2\text{O}}^c \approx P_{22,\text{H}_2\text{O}}^c \approx 1$. A line 4 - $P_A = P^k / P_{\text{H}_2\text{O}}^k$ was drawn by interpolation of values $P_{i,A} = P_i^k / P_{i,\text{H}_2\text{O}}^k$. Values $P_{i,\text{H}_2\text{O}}^k$ were evaluated by $\tau_{i,\text{H}_2\text{O}}^k$ and its spectral dependence (Shukurov 1986), where it was shown, that if w measured in (cm), then $\tau_{i,\text{H}_2\text{O}}^k = \beta \times w^{1.25}$. For $w = 1 \text{ cm}$, numerically, $\beta = \tau_{i,\text{H}_2\text{O}}^k$. Values β for range $v = 765 - 1233 \text{ cm}^{-1}$ are presented in Table 1.

$v_i (\text{cm}^{-1})$	1233	1203	1096	987
i	1	2	10	12
β_i	.038	.040	.045	.053
$v_i (\text{cm}^{-1})$	901	832	765	
i	15	18	22	
β_i	.068	.093	.14	

The aerosol optical depth $\tau_A = -(1/m) \times \ln P_A$ was determined with an accuracy of $\pm 15\%$.

Results

The aerosol optical depth in winter in all “microwindows” i (except $i = 14, 16$) does not exceed $\tau_A = 0.015$. In the summer at large w and at the usual haze conditions, it is more significant, than it was believed. Thus, in $i = 12$, τ_{iA} reaches 0.08. The features of spectral dependence $P_{v,A}$ (its decrease at $\nu \leq 1000 \text{ cm}^{-1}$ and minimum in the range of 900 - 850 cm^{-1}) can be explained by extinction of radiation by water droplets and ice crystals.

Obtained from atmospheric measurement values of k_{i,H_2O}^k (w) agree with laboratory $k_{i,H_2O}^k = \tau_{i,H_2O}^k / w$ of Kulp et al. (1994) within an accuracy of $\pm 10\%$, where τ_{v,H_2O}^k was measured at $e \approx 10.7 \text{ mb}$, $t = +24 \text{ }^\circ\text{C}$. For instance, in $i = 15$ ($\nu = 901 \text{ cm}^{-1}$), $k_{15,H_2O}^k = 0.081 \text{ cm}^{-1}$ by our data and 0.085 cm^{-1} by laboratory data. It confirms the suggestion that the water vapor H_2O continuum is determined practically by, so-called, “e-type”, water vapor absorption of radiation.

The estimates of $P_{i,mol}$ agree with those calculated by Paramonova et al. (1978); however, for $i = 6$, our value $\tau_{6,N_2O} = -(1/m) \times \ln P_{6,N_2O}$ (m), where P_{6,N_2O} (m) = $P_6 / P_{6,N_2O}^o$, more and equal to 0.05 atm^{-1} (compared to calculated data, which was carried out at $m = 4$). The calculated value of τ_{6,N_2O} is equal to 0.028 atm^{-1} by Bignell et al. (1963) and 0.038 atm^{-1} by Paramonova et al. (1978).

In conclusion, it can be noted that the data on solar radiation extinction at small amounts of water vapor content (in middle latitudes - in winter) are very important for analysis of atmospheric transmittance.

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